

PHYS 3313 – Section 001

Lecture # 19

Monday, April 13, 2015

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- Refresher: Infinite Square-well Potential
 - Energy quantization
 - Expectation value computations
- Finite Square Well Potential
- Penetration Depth



Announcements

- Quiz #4 at the beginning of the class this Wednesday, Apr. 15
 - The class web site was down so I will give you an extension..
 - Covers CH 5.4 through what we finish today
 - BFOF
- Colloquium 4pm Wednesday, SH101



Reminder: Special project #5

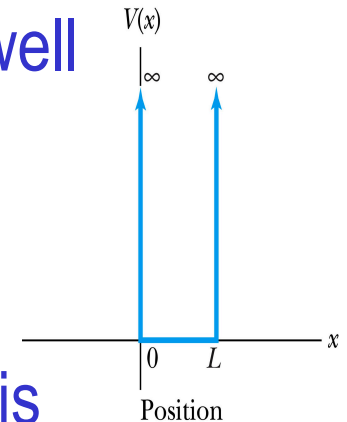
- Show that the Schrodinger equation becomes Newton's second law in the classical limit. (15 points)
- Deadline this Wednesday, Apr. 15, 2015
- You MUST have your own answers!



Infinite Square-Well Potential

- The simplest such system is that of a particle trapped in a box with infinitely hard walls that the particle cannot penetrate. This potential is called an infinite square well and is given by

$$V(x) = \begin{cases} \infty & x \leq 0, x \geq L \\ 0 & 0 < x < L \end{cases}$$



- The wave function must be zero where the potential is infinite.
- Where the potential is zero inside the box, the time independent Schrödinger wave equation $-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + V(x)\psi(x) = E\psi(x)$ becomes $\frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2}\psi = -k^2\psi$ where $k = \sqrt{2mE/\hbar^2}$.
- The general solution is $\psi(x) = A \sin kx + B \cos kx$.

Quantization

- Since the wave function must be continuous, the boundary conditions of the potential dictate that the wave function must be zero at $x = 0$ and $x = L$. These yield valid solutions for $B=0$, and for **integer values** of n such that $kL = n\pi \rightarrow k=n\pi/L$

- The wave function is now
$$\psi_n(x) = A \sin\left(\frac{n\pi x}{L}\right)$$

- We normalize the wave function
$$\int_{-\infty}^{+\infty} \psi_n^*(x) \psi_n(x) dx = 1$$

$$A^2 \int_{-\infty}^{+\infty} \sin^2\left(\frac{n\pi x}{L}\right) dx = A^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx = 1$$

- The normalized wave function becomes

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

- These functions are identical to those obtained for a vibrating string with fixed ends.

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Quantized Energy

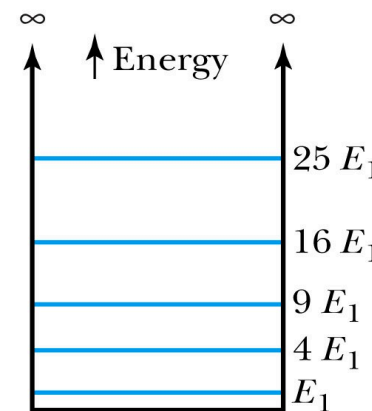
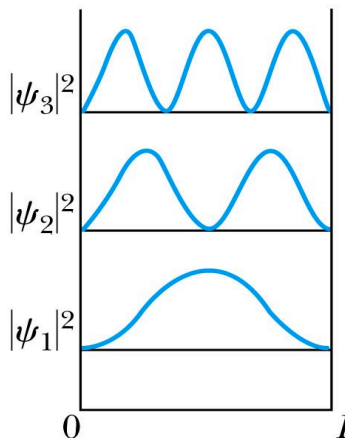
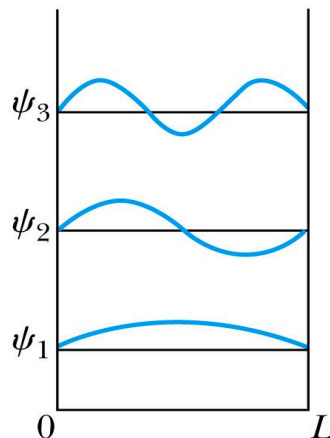
- The quantized wave number now becomes $k_n(x) = \frac{n\pi}{L} = \sqrt{\frac{2mE_n}{\hbar^2}}$
- Solving for the energy yields

$$E_n = n^2 \frac{\pi^2 \hbar^2}{2mL^2} \quad (n = 1, 2, 3, \dots)$$

- Note that the energy depends on the integer values of n . Hence the energy is quantized and nonzero.
- The special case of $n = 1$ is called the **ground state energy**.

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

$$\psi_n^* \psi_n = |\psi_n|^2 = \frac{2}{L} \sin^2\left(\frac{n\pi x}{L}\right)$$



$$E_3 = \frac{9\pi^2 \hbar^2}{2mL^2} = 9E_1$$

$$E_2 = \frac{2\pi^2 \hbar^2}{mL^2} = 4E_1$$

$$E_1 = \frac{\pi^2 \hbar^2}{2mL^2}$$

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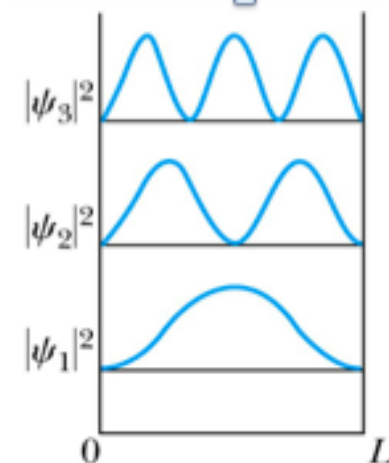
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How does this correspond to Classical Mech.?

- What is the probability of finding a particle in a box of length L ? $\frac{1}{L}$
- Bohr's **correspondence principle** says that QM and CM must correspond to each other! When?
 - When n becomes large, the QM approaches to CM
- So when $n \rightarrow \infty$, the probability of finding a particle in a box of length L is

$$P(x) = \psi_n^*(x)\psi_n(x) = |\psi_n(x)|^2 = \frac{2}{L} \lim_{n \rightarrow \infty} \sin^2\left(\frac{n\pi x}{L}\right) \approx \frac{2}{L} \left\langle \sin^2\left(\frac{n\pi x}{L}\right) \right\rangle = \frac{2}{L} \cdot \frac{1}{2} = \frac{1}{L}$$

- Which is identical to the CM probability!!
- One can also see this from the plot of P !



Expectation Value & Operators

- Expectation value for any function $g(x)$

$$\langle g(x) \rangle = \int_{-\infty}^{+\infty} \Psi(x,t)^* g(x) \Psi(x,t) dx$$

- Position operator is the same as itself, x
- Momentum Operator

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}$$

- Energy Operator

$$\hat{E} = i\hbar \frac{\partial}{\partial t}$$



Ex 6.8: Expectation values inside a box

Determine the expectation values for x , x^2 , p and p^2 of a particle in an infinite square well for the first excited state.

What is the wave function of the first excited state? $n=?$ 2

$$\psi_{n=2}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right)$$

$$\langle x \rangle_{n=2} = \int_{-\infty}^{+\infty} \psi_{n=2}^*(x) x \psi_{n=2}(x) dx = \frac{2}{L} \int_0^L x \sin^2\left(\frac{2\pi x}{L}\right) dx = \frac{L}{2}$$

$$\langle x^2 \rangle_{n=2} = \frac{2}{L} \int_0^L x^2 \sin^2\left(\frac{2\pi x}{L}\right) dx = 0.32L^2$$

$$\langle p \rangle_{n=2} = \frac{2}{L} \int_0^L \sin\left(\frac{2\pi x}{L}\right) (-i\hbar) \frac{\partial}{\partial x} \left[\sin\left(\frac{2\pi x}{L}\right) \right] dx = -i\hbar \frac{2}{L} \frac{2\pi}{L} \int_0^L \sin\left(\frac{2\pi x}{L}\right) \cos\left(\frac{2\pi x}{L}\right) dx = 0$$

$$\langle p^2 \rangle_{n=2} = \frac{2}{L} \int_0^L \sin\left(\frac{2\pi x}{L}\right) (-i\hbar)^2 \frac{\partial^2}{\partial x^2} \left[\sin\left(\frac{2\pi x}{L}\right) \right] dx = \hbar^2 \frac{2}{L} \left(\frac{2\pi}{L}\right)^2 \int_0^L \sin^2\left(\frac{2\pi x}{L}\right) dx = \frac{4\pi^2 \hbar^2}{L^2}$$

$$E_2 = \frac{4\pi^2 \hbar^2}{2mL^2} = \frac{\langle p^2 \rangle_{n=2}}{2m}$$

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Ex 6.9: Proton Transition Energy

A typical diameter of a nucleus is about 10^{-14}m . Use the infinite square-well potential to calculate the transition energy from the first excited state to the ground state for a proton confined to the nucleus.

The energy of the state n is $E_n = n^2 \frac{\pi^2 \hbar^2}{2mL^2}$

What is n for the ground state? $n=1$

$$E_1 = \frac{\pi^2 \hbar^2}{2mL^2} = \frac{\pi^2 \hbar^2 c^2}{2mc^2 L^2} = \frac{1}{mc^2} \frac{\pi^2 \cdot (197.3\text{eV} \cdot \text{nm})^2}{2 \cdot (10^{-5}\text{nm})} = \frac{1.92 \times 10^{15} \text{eV}^2}{938.3 \times 10^6 \text{eV}} = 2.0\text{MeV}$$

What is n for the 1st excited state? $n=2$

$$E_2 = 2^2 \frac{\pi^2 \hbar^2}{2mL^2} = 8.0\text{MeV}$$

So the proton transition energy is

$$\Delta E = E_2 - E_1 = 6.0\text{MeV}$$