$$\frac{d6}{dSl} = \left(\frac{22'e^{2}}{4E}\right)^{2} \frac{1}{Sh^{4}g_{3}}$$

$$= \left(\frac{79x_{2}xe^{2}/f_{5}c}{4 \times 10 \text{ keV}}\right)^{2} \frac{1}{Sh^{4}g_{3}} \frac{e^{2}}{f_{5}c} = 137$$

Note
$$\left(\frac{ZZ^{1}e^{2}}{4e}\right)^{2} = \left(\frac{2\times79\times(4.8\times10^{-10}e^{2})^{2}}{4\times1810^{4}e^{2}\times10^{10}e^{2}}\right)^{2}$$

 $erg = \frac{9\cdot cm^{2}}{5^{2}}$ $esu = \frac{9^{1}\pi cm^{3/2}}{c^{2}} = \left(5.7\times10^{-10}cm\right)^{2}$

=271 (ZZ'e2)2 So Sno 1 suren do

2.) Gror = Sanda Sitt do (0) snodo

= -4 17 × 3,25×105 Barry & (-14) = 5,71×107 Barry 6=300 (-3) = 1,22×107 Barry 6=600

(1-2) 0=90

(-1) = 40x106 Bing 6=90°

Problem 1.10 What is the minimum impact parameter needed to deflect 7.7 MeV α -particles from gold nuclei by at least 1°? What about by at least 30°? What is the ratio of probabilities for deflections of $\theta > 1^{\circ}$ relative to $\theta > 30^{\circ}$? (See the CRC Handbook for the density of gold.)

For the scattering of a 7.7 MeV α -particle from gold, we have

$$Z = 2$$
, $Z' = 79$, $E = 7.7 \,\text{MeV}$, (1.74)

so that we obtain

$$\begin{split} \frac{ZZ'e^2}{2E} &= ZZ' \times \frac{\hbar c}{2E} \times \frac{e^2}{\hbar c} \\ &= 2 \times 79 \times \frac{197 \,\text{MeV} - \text{F}}{2 \times 7.7 \,\text{MeV}} \times \frac{1}{137} \\ &\approx 14.5 \times 10^{-13} \,\text{cm} \approx 1.4 \times 10^{-12} \,\text{cm}. \end{split} \tag{1.75}$$

We know from Eq. (1.32) of the text that

$$b = \frac{ZZ'e^2}{2E}\cot\frac{\theta}{2}, \qquad (1.76)$$

which leads to

$$b(\theta) \approx 1.4 \times 10^{-12} \cot \frac{\theta}{2} \text{ cm.}$$
 (1.77)

We note that for

$$\theta = 1^{\circ} = \frac{\pi}{180} \approx \frac{1}{60} \ll 1,$$
 (1.78)

we have

$$\tan\frac{\theta}{2} \approx \frac{\theta}{2} \approx \frac{1}{120}, \quad \cot\frac{\theta}{2} = \frac{1}{\tan\frac{\theta}{5}} \approx 120.$$
 (1.79)

Similarly, for

$$\theta = 30^{\circ} = \frac{\pi}{6} \approx \frac{1}{2} \qquad (1.80)$$

we have

$$\tan \frac{\theta}{2} \approx \frac{\theta}{2} \approx \frac{1}{4}, \quad \cot \frac{\theta}{2} = \frac{1}{\tan \frac{\theta}{2}} \approx 4.$$
 (1.81)

Using these values in (1.77), we obtain

$$b(\theta = 1^{\circ}) \approx 1.4 \times 10^{-12} \times 120 \,\mathrm{cm} = 1.7 \times 10^{-10} \,\mathrm{cm},$$

 $b(\theta = 30^{\circ}) \approx 1.4 \times 10^{-12} \times 4 \,\mathrm{cm} = 5.6 \times 10^{-12} \,\mathrm{cm}.$ (1.82)

As we have already seen in Problem 1.1, the probability of scattering for angles greater than θ_b goes as the area πb^2 . Therefore, using (1.82) we have

$$\frac{\sigma(\theta > 1^{\circ})}{\sigma(\theta > 30^{\circ})} = \frac{b^{2}(\theta = 1^{\circ})}{b^{2}(\theta = 30^{\circ})}$$

$$\approx \left(\frac{1.7 \times 10^{-10}}{5.6 \times 10^{-12}}\right)^{2} \approx 900. \quad (1.83)$$

In other words, there will be approximately 900 more particle collisions for $\theta > 1^{\circ}$ than for $\theta > 30^{\circ}$. 20 Solutions Manual

we have

$$\tan\frac{\theta}{2} \approx \frac{\theta}{2} \approx \frac{1}{120}, \quad \cot\frac{\theta}{2} = \frac{1}{\tan\frac{\theta}{2}} \approx 120.$$
(1.79)

Similarly, for

$$\theta = 30^{\circ} = \frac{\pi}{6} \approx \frac{1}{2} \tag{1.80}$$

we have

$$\tan\frac{\theta}{2} \approx \frac{\theta}{2} \approx \frac{1}{4}, \quad \cot\frac{\theta}{2} = \frac{1}{\tan\frac{\theta}{2}} \approx 4.$$
(1.81)

Using these values in (1.77), we obtain

$$b(\theta = 1^{\circ}) \approx 1.4 \times 10^{-12} \times 120 \,\mathrm{cm} = 1.7 \times 10^{-10} \,\mathrm{cm},$$

 $b(\theta = 30^{\circ}) \approx 1.4 \times 10^{-12} \times 4 \,\mathrm{cm} = 5.6 \times 10^{-12} \,\mathrm{cm}.$ (1.82)

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$$\approx \left(\frac{1.7 \times 10^{-10}}{5.6 \times 10^{-12}}\right)^{2} \approx 900.$$
(1.83)

In other words, there will be approximately 900 more particle collisions for $\theta > 1^{\circ}$ than for $\theta > 30^{\circ}$.

Problem 1.11 Consider a collimated source of 8 MeV α -particles that provides 10^4 α/sec that impinge on a 0.1 mm gold foil. What counting rate would you expect in a detector that subtends an annular cone of $\Delta\theta=0.05$ rad, at a scattering angle of $\theta=90^\circ$? Compare this to the rate at $\theta=5^\circ$. Is there a problem? Is it serious (see Problem 1.12). (Hint: You can use the small-angle approximation where appropriate, and find the density of gold in the CRC Handbook.)

$$Z=2, \quad Z'=79, \quad E=8\,\mathrm{MeV},$$
 $\rho=\mathrm{density} \ \mathrm{of} \ \mathrm{gold} \approx 19.3\,\mathrm{g/cm^3},$
 $t=\mathrm{Thickness} \ \mathrm{of} \ \mathrm{gold} \ \mathrm{foil}=0.1\,\mathrm{mm}=10^{-2}\,\mathrm{cm},$
 $N_0=\mathrm{Incident} \ \mathrm{flux}=10^4/\mathrm{sec},$
 $A=\mathrm{Atomic} \ \mathrm{weight} \ \mathrm{of} \ \mathrm{gold}=197,$

$$(1.84)$$

 $A_0 = \text{Avogadro's number} \approx 6 \times 10^{23}/\text{mole},$

 $\Delta \theta = \text{Angle subtended by the detector} = 0.05 \, \text{rad}.$

We can therefore calculate

$$\frac{d\sigma}{d\Omega}(\theta) = \left(\frac{ZZ'e^2}{4E}\right)^2 \csc^4 \frac{\theta}{2} (sr)^{-1}$$

$$= \left(ZZ' \times \frac{\hbar c}{4E} \times \frac{e^2}{\hbar c}\right)^2 \csc^4 \frac{\theta}{2} (sr)^{-1}$$

$$= \left(2 \times 79 \times \frac{197 \text{ MeV} - F}{4 \times 8 \text{ MeV}} \times \frac{1}{137}\right)^2 \csc^4 \frac{\theta}{2} (sr)^{-1}$$

$$\approx 0.48 \times 10^{-24} \csc^4 \frac{\theta}{2} \text{ cm}^2/\text{sr}.$$
(1.85)

Similarly, we have

$$\frac{A_0 \rho t}{A} \approx \frac{6 \times 10^{23} \times 19.3 \times 10^{-2} / \text{cm}^2}{197}$$
$$\approx 6 \times 10^{20} / \text{cm}^2. \tag{1.86}$$

From Eq. (1.40) of the text we therefore obtain the counting rate $\mathrm{d}n(\theta)=N_0\frac{A_0\rho t}{A}\,\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega}(\theta)\,\mathrm{d}\Omega$

$$\approx 10^4/\text{sec} \times 6 \times 10^{20}/\text{cm}^2 \times 0.48 \times 10^{-24} \text{cosec}^4 \frac{\theta}{2} \text{cm}^2/\text{sr} \times d\Omega$$
$$= 2.88 \operatorname{cosec}^4 \frac{\theta}{2} d\Omega (\text{sec} - \text{sr})^{-1}. \tag{1.87}$$

For scattering with azimuthal symmetry, we can write

$$\mathrm{d}\Omega = 2\pi\sin\theta\mathrm{d}\theta,\tag{1.88}$$

and if we identify $d\theta \approx \Delta\theta = 0.05 \, \mathrm{rad}$, we get

$$d\Omega \approx 2\pi \sin \theta \times 0.05 \,\text{sr} \approx 0.3 \sin \theta \,\text{sr}.$$
 (1.89)

Putting this back into (1.87), we obtain

$$dn(\theta) \approx 2.88 \operatorname{cosec}^{4} \frac{\theta}{2} (\operatorname{sec} - \operatorname{sr})^{-1} \times 0.3 \sin \theta \operatorname{sr}$$
$$\approx 0.86 \sin \theta \operatorname{cosec}^{4} \frac{\theta}{2} (\operatorname{sec})^{-1}. \tag{1.90}$$

It follows that

$$dn\left(\theta = \frac{\pi}{2}\right) \approx 0.86 \times 1 \times (\sqrt{2})^4 \approx 3.4/\text{sec.}$$
 (1.91)

On the other hand, for $\theta = 5^{\circ} = \frac{\pi}{36} \approx \frac{1}{12} \ll 1$, we have

$$\sin \theta \approx \theta \approx \frac{1}{12}, \quad \csc^4 \frac{\theta}{2} \approx \left(\frac{2}{\theta}\right)^4 \approx (24)^4, \quad (1.92)$$

and we obtain

$$dn(\theta = 5^{\circ}) \approx 0.86 \times \frac{1}{12} \times (24)^{4}/\text{sec} \approx 2.4 \times 10^{4}/\text{sec}.$$
 (1.93)

This is, in fact, larger than the incident flux, and, if this were true, conservation of probability (particle number) would be violated, which is a serious problem! For one thing, we note that the approximation

$$d\theta \approx \Delta\theta,$$
 (1.94)

is meaningful only when

$$\frac{\Delta\theta}{\theta} \ll 1,\tag{1.95}$$

which is clearly violated when $\Delta\theta=0.05$ rad and $\theta=5^{\circ}\approx\frac{1}{12}\approx0.08$ rad. This is one of the sources of the difficulty. For other sources of this error, we turn to the solution of the next problem.

Problem 1.12 Consider the expression Eq. (1.41) for Rutherford Scattering of α -particles from gold nuclei. Integrate this over all angles to obtain n. In principle, n cannot exceed N_0 , the number of incident particles. Why? What cutoff value for θ would be required in the integral, that is, some $\theta = \theta_0 > 0$, to assure that n does not exceed N_0 in Problem 1.4? (Hint: After integrating, use the small-angle approximation to simplify the calculation.) Using the Heisenberg uncertainty principle $\Delta p_x \Delta x \approx \hbar$, where Δx is some