

# PHYS 3446 – Lecture #3

*Tuesday, Jan. 27 2015*

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1. Rutherford Scattering with Coulomb force
2. Scattering Cross Section
3. Differential Cross Section of Rutherford Scattering
4. Measurement of Cross Sections



# Elastic Scattering



- From momentum and kinetic energy conservation

$$v_0^2 = v_\alpha^2 + \frac{m_t}{m_\alpha} v_t^2 \quad \text{***Eq. 1.2}$$

$$v_t^2 \left( 1 - \frac{m_t}{m_\alpha} \right) = 2\vec{v}_\alpha \cdot \vec{v}_t \quad \text{***Eq. 1.3}$$



## Analysis Case 1

- If  $m_t \ll m_\alpha$ ,

– change of momentum of alpha particle is negligible

$$v_t^2 \left( 1 - \frac{m_t}{m_\alpha} \right) = 2\vec{v}_\alpha \cdot \vec{v}_t$$

## Analysis Case 2

- If  $m_t \gg m_\alpha$ ,

– alpha particle deflected backwards

# Rutherford Scattering with EM Force 1

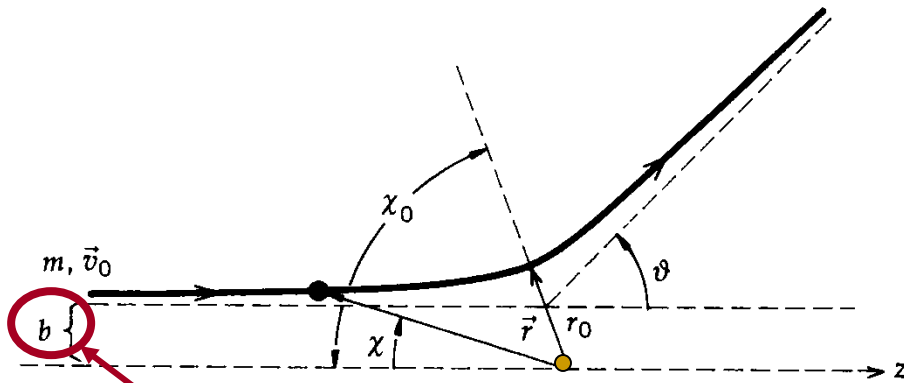


- Need to take into account the EM force between the  $\alpha$  and the atom
- The Coulomb force is a central force  $\rightarrow$  conservative force
- The Coulomb potential energy between particles with  $Ze$  and  $Z'e$  electrical charge, separated by a distance  $r$  is
$$V(r) = \frac{ZZ'e^2}{r}$$
- Since the total energy is conserved,

$$E = \frac{1}{2}mv_0^2 = \text{constant} > 0 \quad \Rightarrow \quad v_0 = \sqrt{\frac{2E}{m}}$$



# Rutherford Scattering with EM Force 2



Impact parameter=distance of closest approach if there were no Coulomb force

- The distance vector  $\mathbf{r}$  is always the same direction as the force throughout the entire motion, so the net torque ( $\mathbf{r} \times \mathbf{F}$ ) is 0.
- Since there is no net torque, the angular momentum ( $\mathbf{l} = \mathbf{r} \times \mathbf{p}$ ) is conserved.  $\rightarrow$  The magnitude of the angular momentum is  $l = mvb$ .

- From the energy relation, we obtain

$$l = m\sqrt{2E/m} \cdot b = b\sqrt{2mE} \Rightarrow b^2 = l^2/2mE$$

- From angular momentum defn. , we obtain an equation of motion , where  $\chi =$  angle between  $-x$  axis and radial vector"  $d\chi/dt = l/mr^2$
- From energy conservation, we obtain another equation of motion

$$E = \frac{1}{2}m\left(\frac{dr}{dt}\right)^2 + \frac{1}{2}mr^2\left(\frac{d\chi}{dt}\right)^2 + V(r)$$

$$\Rightarrow \frac{dr}{dt} = \pm \sqrt{\frac{2}{m} \left( E - V(r) - \frac{l^2}{2mr^2} \right)}$$

Centrifugal barrier

Effective potential



# Rutherford Scattering with EM Force 3

- Rearranging the terms, we obtain

Eq. 1.18

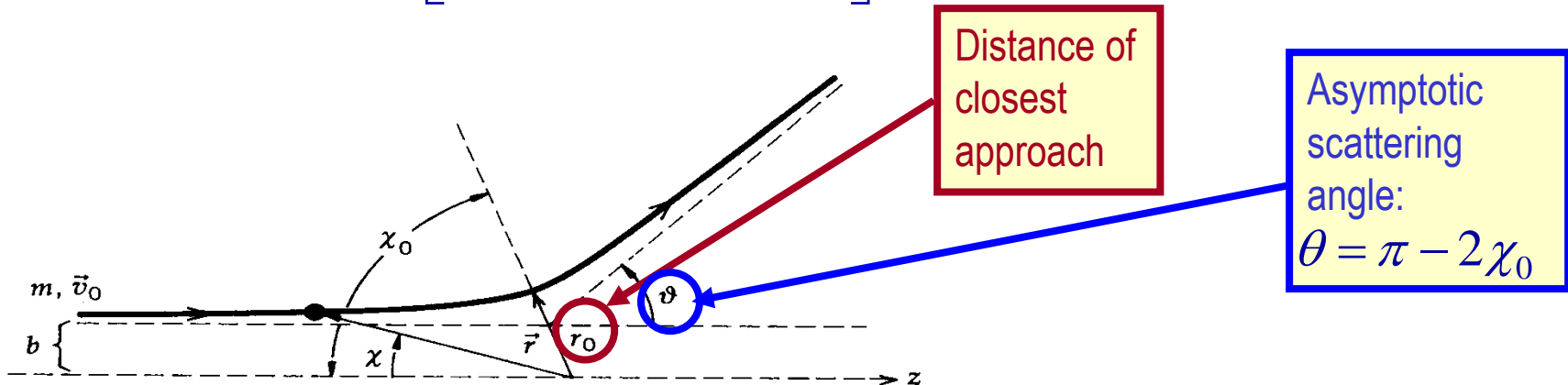
$$\frac{dr}{dt} = -\frac{l}{mrb} \sqrt{r^2 \left( 1 - \frac{V(r)}{E} \right) - b^2}$$

-sign since r decreases as alpha particle approaches target

Using 1.18 in the following  $d\chi/dt = L/mr^2 = d\chi/dr \cdot dr/dt$  gives

- and

$$d\chi = -\frac{bdr}{r \left[ r^2 \left( 1 - \frac{V(r)}{E} \right) - b^2 \right]^{1/2}} \quad \text{Eq. 1.19}$$





# Rutherford Scattering with EM Force 4

- What happens at the DCA?
  - Kinetic energy goes to 0.  $\left. \frac{dr}{dt} \right|_{r=r_0} = 0$
  - Consider the case where the alpha particle is incident on the z-axis, it would reach the DCA, stop, and reverse direction!
  - From Eq. 1.18, we can obtain  $r_0^2 \left( 1 - \frac{V(r_0)}{E} \right) - b^2 = 0$
  - This allows us to determine DCA for a given potential and  $\chi_0$ .
- Substituting 1.19 into the scattering angle equation gives:

$$\theta = \pi - 2\chi_0 = \pi - 2b \int_{r_0}^{\infty} \frac{dr}{r \left[ r^2 \left( 1 - \frac{V(r)}{E} \right) - b^2 \right]^{1/2}}$$



# Rutherford Scattering with EM Force 5

- For a Coulomb potential

$$V(r) = \frac{ZZ'e^2}{r}$$

- DCA can be obtained for a given impact parameter  $b$ ,

$$r_0 = \frac{ZZ'e^2/E}{2} \left( 1 + \sqrt{1 + 4b^2 E^2 / (ZZ'e^2)^2} \right)$$

- And the angular distribution becomes

$$\theta = \pi - 2b \int_{r_0}^{\infty} \frac{dr}{r \left[ r^2 \left( 1 - \frac{ZZ'e^2}{rE} \right) - b^2 \right]^{1/2}}$$





# Rutherford Scattering with EM Force 6

- Replace the variable  $1/r=x$ , and performing the integration, we obtain

$$\theta = \pi + 2b \cos^{-1} \left( \frac{1}{\sqrt{1 + 4b^2 E^2 / (ZZ' e^2)^2}} \right)$$

- This can be rewritten

$$\frac{1}{\sqrt{1 + 4b^2 E^2 / (ZZ' e^2)^2}} = \cos \left( \frac{\theta - \pi}{2} \right)$$

- Solving this for  $b$ , we obtain

$$b = \frac{ZZ' e^2}{2E} \cot \frac{\theta}{2}$$

# Rutherford Scattering with EM Force 7



$$b = \frac{ZZ'e^2}{2E} \cot \frac{\theta}{2}$$

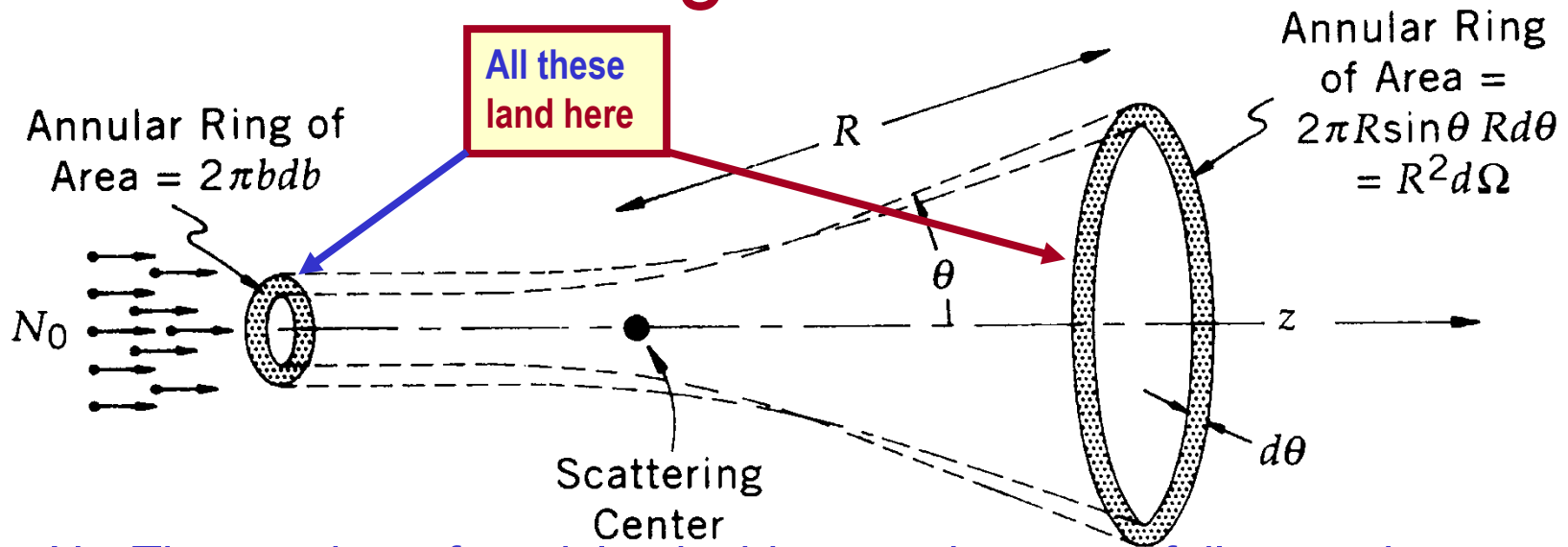
- From the solution for  $b$ , we can learn the following
  1. For fixed  $b$  and  $E$ 
    - The scattering is larger for a larger value of  $Z$  or  $Z'$  (large charge in projectile or target)
      - Makes perfect sense since Coulomb potential is stronger with larger  $Z$ .
      - Results in larger deflection.
  2. For fixed  $b$ ,  $Z$  and  $Z'$ 
    - The scattering angle is larger when  $E$  is smaller.
      - If particle has low energy, its velocity is smaller
      - Spends more time in the potential, suffering greater deflection
  3. For fixed  $Z$ ,  $Z'$ , and  $E$ 
    - The scattering angle is larger for smaller impact parameter  $b$ 
      - Makes perfect sense also, since as the incident particle is closer to the nucleus, it feels stronger Coulomb force.

# What do we learn from scattering?



- Scattering of a particle in a potential is completely determined when we know both
    - The impact parameter,  $b$ , and
    - The energy of the incident particle,  $E$
- $$b = \frac{ZZ'e^2}{2E} \cot \frac{\theta}{2}$$
- For a fixed energy, the deflection is defined by
    - The impact parameter,  $b$ .
  - What do we need to perform a scattering experiment?
    - Incident flux of beam particles with known  $E$
    - Device that can measure number of scattered particles at various angle,  $\theta$ .
    - Measurements of the number of scattered particles reflect
      - Impact parameters of the incident particles
      - The effective size of the scattering center
  - By measuring the scattering angle  $\theta$ , we can learn about the potential or the forces between the target and the projectile

# Scattering Cross Section



- $N_0$ : The number of particles incident on the target foil per unit area per unit time.
- Any incident particles entering with impact parameter  $b$  and  $b+db$  will scatter to the angle  $\theta$  and  $\theta-d\theta$ .
- In other words, they scatter into the solid angle  $d\Omega$  ( $=2\pi\sin\theta d\theta$ ).
- So the number of particles scattered into the solid angle  $d\Omega$  per unit time is  $2\pi N_0 bdb$ .
- Note: have assumed thin foil, and large separation between nuclei—why?

# Scattering Cross Section



- For a central potential
  - Such as Coulomb potential
  - Which has spherical symmetry
- The scattering center presents an effective transverse cross-sectional area of

$$\Delta\sigma = 2\pi b db$$

- For the particles to scatter into  $\theta$  and  $\theta+d\theta$



# Scattering Cross Section

- In more generalized cases,  $\Delta\sigma$  depends on both  $\theta$  &  $\phi$ .

$$\Delta\sigma(\theta, \phi) = b db d\phi = - \frac{d\sigma}{d\Omega}(\theta, \phi) d\Omega = - \frac{d\sigma}{d\Omega}(\theta, \phi) \sin\theta d\theta d\phi$$

Why negative?

Since the deflection and change of b are in opposite direction!!

- With a spherical symmetry,  $\phi$  can be integrated out:

$$\Delta\sigma(\theta) = - \frac{d\sigma}{d\Omega}(\theta) 2\pi \sin\theta d\theta = 2\pi b db$$

reorganize

$$\frac{d\sigma}{d\Omega}(\theta) = - \frac{b}{\sin\theta} \frac{db}{d\theta}$$

Differential Cross Section

What is the dimension of the differential cross section?

Area!!

# Scattering Cross Section

femptobarn



- For a central potential, measuring the yield as a function of  $\theta$  (the differential cross section) is equivalent to measuring the entire effect of the scattering
- So what is the physical meaning of the differential cross section?

⇒ Measurement of yield as a function of specific experimental variables

⇒ This is equivalent to measuring the probability of occurrence of a physical process in a specific kinematic phase space

- Cross sections are measured in the unit of barns:

$$1 \text{ barn} \equiv 10^{-24} \text{ cm}^2$$

picobarn

Where does this come from?

Cross sectional area of a typical nucleus!

nanobarn

15

# Cross Section of Rutherford Scattering

- The impact parameter in Rutherford scattering is

$$b = \frac{ZZ'e^2}{2E} \cot \frac{\theta}{2}$$

- Thus,

$$\frac{db}{d\theta} = -\frac{1}{2} \frac{ZZ'e^2}{2E} \operatorname{cosec}^2 \frac{\theta}{2}$$

- Differential cross section of Rutherford scattering is

$$\frac{d\sigma}{d\Omega}(\theta) = -\frac{b}{\sin \theta} \frac{db}{d\theta} = \left( \frac{ZZ'e^2}{4E} \right)^2 \operatorname{cosec}^4 \frac{\theta}{2} = \left( \frac{ZZ'e^2}{4E} \right)^2 \frac{1}{\sin^4 \frac{\theta}{2}}$$

what happened? can you say “trig identity?”  $\sin(2x) = 2\sin(x) \cos(x)$



# Measuring Cross Sections

The diagram shows the equation  $dn = \frac{N_0 \rho t A_0}{A} \frac{d\sigma(\theta, \phi)}{d\Omega} d\Omega = N_0 \frac{N}{S} \frac{d\sigma(\theta, \phi)}{d\Omega} d\Omega$ . The terms are annotated with boxes and arrows:

- $dn$  is circled in red, with an arrow pointing to a box labeled "Number of detected particles/sec".
- $N_0$  is circled in red, with an arrow pointing to a box labeled "Projectile particle flux".
- $N$  is circled in red, with an arrow pointing to a box labeled "Density of the target particles".
- $\frac{d\sigma(\theta, \phi)}{d\Omega}$  is circled in red, with an arrow pointing to a box labeled "Scattering cross section".
- $d\Omega$  is circled in red, with an arrow pointing to a box labeled "Detector acceptance".

- This is a general expression for any scattering process, independent of the theory
- This gives an observed counts per second