

# PHYS 1443 – Section 003

## Lecture #14

*Monday, Nov. 4, 2002*

*Dr. Jaehoon Yu*

1. Parallel Axis Theorem
2. Torque
3. Torque & Angular Acceleration
4. Work, Power and Energy in Rotation

Today's homework is homework #14 due 12:00pm, Monday, Nov. 11!!

Monday, Nov. 4, 2002



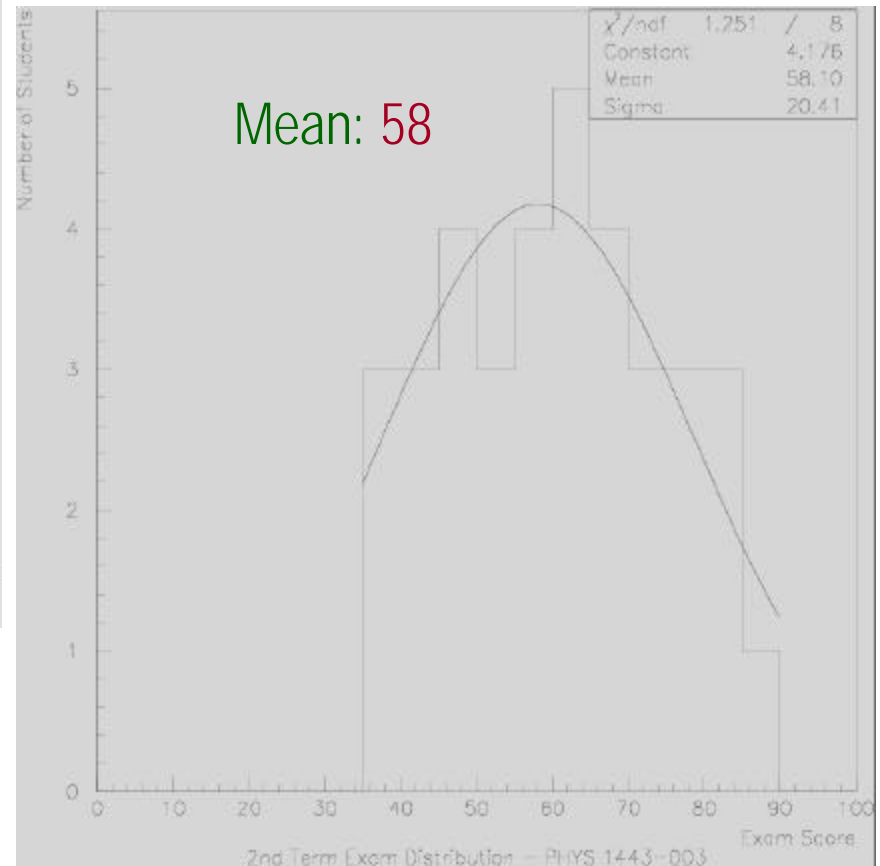
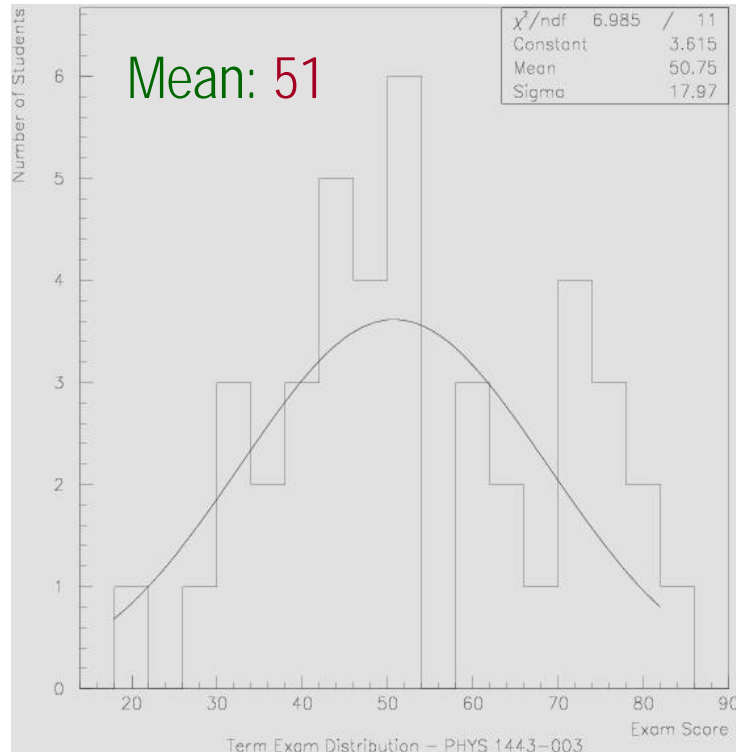
PHYS 1443-003, Fall 2002  
Dr. Jaehoon Yu

# Announcements

- 2<sup>nd</sup> Term Exam
  - Grading is completed
    - Maximum Score: 87
    - Numerical Average: 58.1
    - Four persons missed the exam without a prior approval
  - Can look at your exam after the class
  - All scores are relative based on the curve
    - One worst after the adjustment will be dropped
  - Exam constitutes only 50% of the total
    - Do your homework well
    - Come to the class and do well with quizzes



# 2<sup>nd</sup> Term Exam Distributions



14% Improvement

A lot narrower distribution. → Even improvements

But as always, you could do better!!!

Monday, Nov. 4, 2002



PHYS 1443-003, Fall 2002  
Dr. Jaehoon Yu

# Calculation of Moments of Inertia

Moments of inertia for large objects can be computed, if we assume the object consists of small volume elements with mass,  $\Delta m_i$ .

The moment of inertia for the large rigid object is

$$I = \lim_{\Delta m_i \rightarrow 0} \sum_i r_i^2 \Delta m_i = \int r^2 dm$$

It is sometimes easier to compute moments of inertia in terms of volume of the elements rather than their mass

How can we do this?

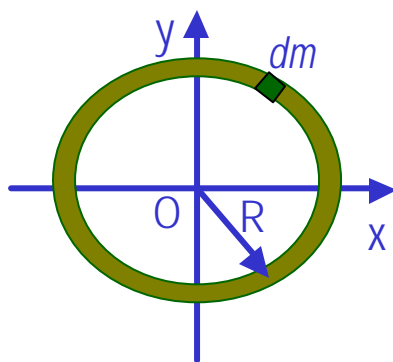
Using the volume density,  $\rho$ , replace  $dm$  in the above equation with  $dV$ .

$$\rho = \frac{dm}{dV} \quad dm = \rho dV$$

The moments of inertia becomes

$$I = \int \rho r^2 dV$$

Example 10.5: Find the moment of inertia of a uniform hoop of mass  $M$  and radius  $R$  about an axis perpendicular to the plane of the hoop and passing through its center.



The moment of inertia is

$$I = \int r^2 dm = R^2 \int dm = MR^2$$

What do you notice from this result?

The moment of inertia for this object is the same as that of a point of mass  $M$  at the distance  $R$ .

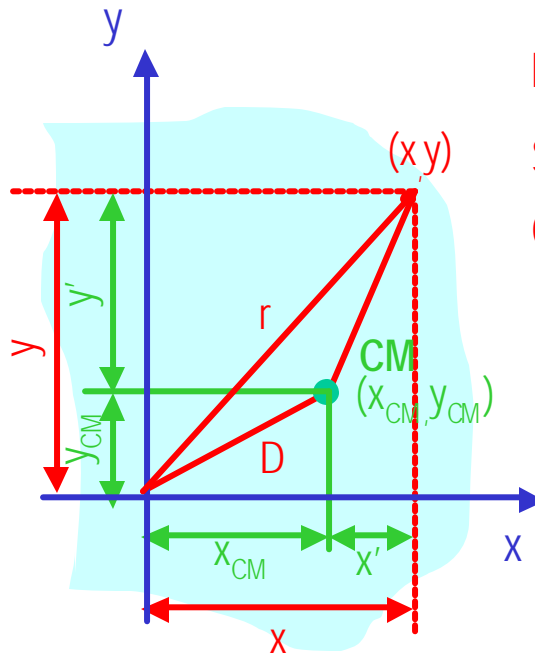
Monday, Nov. 4, 2002



PHYS 1443-003, Fall 2002  
Dr. Jaehoon Yu

# Parallel Axis Theorem

Moments of inertia for highly symmetric object is easy to compute if the rotational axis is the same as the axis of symmetry. However if the axis of rotation does not coincide with axis of symmetry, the calculation can still be done in simple manner using **parallel-axis theorem**.  $I = I_{CM} + MD^2$



Moment of inertia is defined  $I = \int r^2 dm = \int \sqrt{(x^2 + y^2)} dm$  (1)

Since  $x$  and  $y$  are  $x = x_{CM} + x'$   $y = y_{CM} + y'$

One can substitute  $x$  and  $y$  in Eq. 1 to obtain

$$I = \int [(x_{CM} + x')^2 + (y_{CM} + y')^2] dm$$

$$= (x_{CM}^2 + y_{CM}^2) \int dm + 2x_{CM} \int x' dm + 2y_{CM} \int y' dm + \int (x'^2 + y'^2) dm$$

Since the  $x'$  and  $y'$  are the distance from CM, by definition

$$\int x' dm = 0 \quad \int y' dm = 0$$

Therefore, the parallel-axis theorem

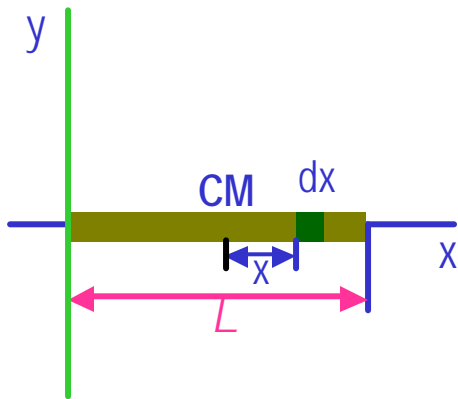
$$I = (x_{CM}^2 + y_{CM}^2) \int dm + \int (x'^2 + y'^2) dm = MD^2 + I_{CM}$$

What does this theorem tell you?

Moment of inertia of any object about any arbitrary axis are the same as the sum of moment of inertia for a rotation about the CM and that of the CM about the rotation axis.

# Example 10.8

Calculate the moment of inertia of a uniform rigid rod of length  $L$  and mass  $M$  about an axis that goes through one end of the rod, using parallel-axis theorem.



The line density of the rod is

$$I = \frac{M}{L}$$

so the masslet is

$$dm = I dx = \frac{M}{L} dx$$

The moment of inertia about the CM

$$\begin{aligned} I_{CM} &= \int r^2 dm = \int_{-L/2}^{L/2} \frac{x^2 M}{L} dx = \frac{M}{L} \left[ \frac{1}{3} x^3 \right]_{-L/2}^{L/2} \\ &= \frac{M}{3L} \left[ \left( \frac{L}{2} \right)^3 - \left( -\frac{L}{2} \right)^3 \right] = \frac{M}{3L} \left( \frac{L^3}{4} \right) = \frac{ML^2}{12} \end{aligned}$$

Using the parallel axis theorem

$$I = I_{CM} + D^2 M = \frac{ML^2}{12} + \left( \frac{L}{2} \right)^2 M = \frac{ML^2}{12} + \frac{ML^2}{4} = \frac{ML^2}{3}$$

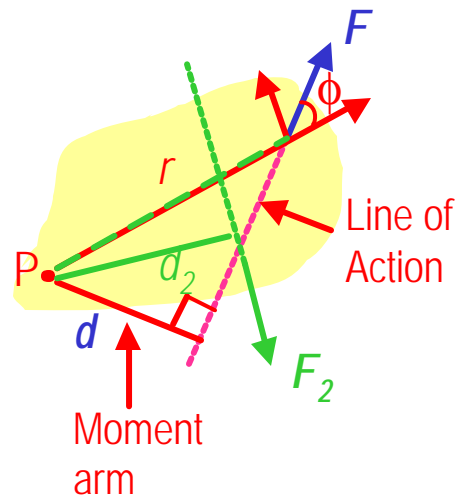
The result is the same as using the definition of moment of inertia.

Parallel-axis theorem is useful to compute moment of inertia of a rotation of a rigid object with complicated shape about an arbitrary axis



# Torque

Torque is the tendency of a force to rotate an object about an axis.  
Torque,  $\tau$ , is a vector quantity.



Consider an object pivoting about the point  $P$  by the force  $F$  being exerted at a distance  $r$ .

The line that extends out of the tail of the force vector is called the **line of action**.

The perpendicular distance from the pivoting point  $P$  to the line of action is called **Moment arm**.

Magnitude of torque is defined as the product of the force exerted on the object to rotate it and the moment arm.

$$\tau \equiv rF \sin \phi = Fd$$

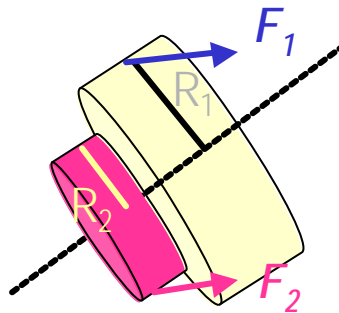
When there are more than one force being exerted on certain points of the object, one can sum up the torque generated by each force vectorially. The convention for sign of the torque is positive if rotation is in counter-clockwise and negative if clockwise.

$$\begin{aligned} \sum \tau &= \tau_1 + \tau_2 \\ &= Fd - F_2 d_2 \end{aligned}$$



## Example 10.9

A one piece cylinder is shaped as in the figure with core section protruding from the larger drum. The cylinder is free to rotate around the central axis shown in the picture. A rope wrapped around the drum whose radius is  $R_1$  exerts force  $F_1$  to the right on the cylinder, and another force exerts  $F_2$  on the core whose radius is  $R_2$  downward on the cylinder. A) What is the net torque acting on the cylinder about the rotation axis?



The torque due to  $F_1$   $t_1 = -R_1 F_1$  and due to  $F_2$   $t_2 = R_2 F_2$

So the total torque acting on the system by the forces is

$$\sum t = t_1 + t_2 = -R_1 F_1 + R_2 F_2$$

Suppose  $F_1 = 5.0 \text{ N}$ ,  $R_1 = 1.0 \text{ m}$ ,  $F_2 = 15.0 \text{ N}$ , and  $R_2 = 0.50 \text{ m}$ . What is the net torque about the rotation axis and which way does the cylinder rotate from the rest?

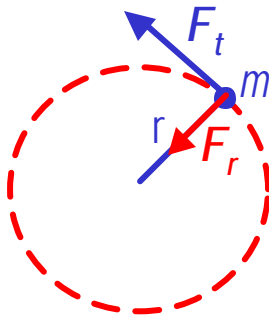
Using the above result

$$\begin{aligned} \sum t &= -R_1 F_1 + R_2 F_2 \\ &= -5.0 \times 1.0 + 15.0 \times 0.50 = 2.5 \text{ N} \cdot \text{m} \end{aligned}$$

The cylinder rotates in counter-clockwise.



# Torque & Angular Acceleration



Let's consider a point object with mass  $m$  rotating on a circle.

What forces do you see in this motion?

The tangential force  $F_t$  and radial force  $F_r$

The tangential force  $F_t$  is

$$F_t = ma_t = mra$$

The torque due to tangential force  $F_t$  is

$$\tau = F_t r = ma_t r = mr^2 a$$

What do you see from the above relationship?

$$\tau = Ia$$

What does this mean?

Torque acting on a particle is proportional to the angular acceleration.

What law do you see from this relationship?

Analogous to Newton's 2<sup>nd</sup> law of motion in rotation.

How about a rigid object?

The external tangential force  $dF_t$  is

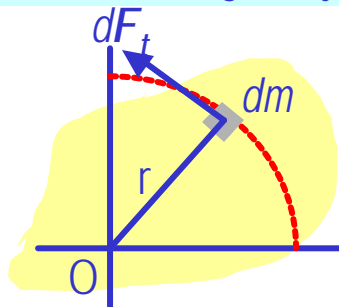
$$dF_t = dma_t = dmra$$

The torque due to tangential force  $F_t$  is

$$d\tau = dF_t r = (r^2 dm)a$$

The total torque is

$$\sum \tau = a \int r^2 dm = Ia$$



Monday, Nov. 4, 2002

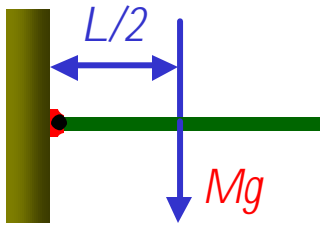
What is the contribution due to radial force and why?

Contribution from radial force is 0, because its line of action passes through the pivoting point, making the moment arm 0.

002

## Example 10.10

A uniform rod of length  $L$  and mass  $M$  is attached at one end to a frictionless pivot and is free to rotate about the pivot in the vertical plane. The rod is released from rest in the horizontal position. What are the initial angular acceleration of the rod and the initial linear acceleration of its right end?



The only force generating torque is the gravitational force  $Mg$

$$\tau = Fd = F \frac{L}{2} = Mg \frac{L}{2} = I\alpha$$

Since the moment of inertia of the rod when it rotates about one end

$$I = \int_0^L r^2 dm = \int_0^L x^2 \frac{M}{L} dx = \left( \frac{M}{L} \right) \left[ \frac{x^3}{3} \right]_0^L = \frac{ML^2}{3}$$

We obtain

$$\alpha = \frac{MgL}{2I} = \frac{MgL}{2 \frac{ML^2}{3}} = \frac{3g}{2L}$$

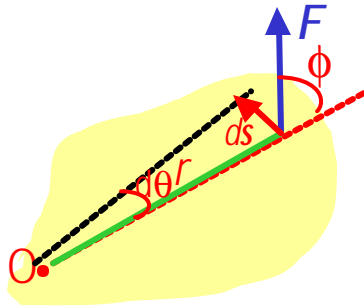
Using the relationship between tangential and angular acceleration

$$a_t = L\alpha = \frac{3g}{2}$$

What does this mean?

The tip of the rod falls faster than an object undergoing a free fall.

# Work, Power, and Energy in Rotation



Let's consider a motion of a rigid body with a single external force  $F$  exerting on the point P, moving the object by  $ds$ .

The work done by the force  $F$  as the object rotates through the infinitesimal distance  $ds=r d\theta$  is

$$dW = \vec{F} \cdot d\vec{s} = (F \sin \phi) r d\theta$$

What is  $F \sin \phi$ ?

The tangential component of force  $F$ .

What is the work done by radial component  $F \cos \phi$ ?

Zero, because it is perpendicular to the displacement.

Since the magnitude of torque is  $r F \sin \phi$ ,

$$dW = \tau d\theta$$

The rate of work, or power becomes

$$P = \frac{dW}{dt} = \frac{\tau d\theta}{dt} = \tau \omega$$

How was the power defined in linear motion?

The rotational work done by an external force equals the change in rotational energy.

$$\sum \tau = I \alpha = I \left( \frac{d\omega}{dt} \right) = I \left( \frac{d\omega}{d\theta} \right) \left( \frac{d\theta}{dt} \right)$$

The work put in by the external force then

$$dW = \sum \tau d\theta = I \omega d\omega$$

$$W = \int_{\theta_i}^{\theta_f} \sum \tau d\theta = \int_{\omega_i}^{\omega_f} I \omega d\omega = \frac{1}{2} I \omega_f^2 - \frac{1}{2} I \omega_i^2$$

Monday, Nov. 4, 2002



PHYS 1443-003, Fall 2002  
Dr. Jaehoon Yu

# Similarity Between Linear and Rotational Motions

All physical quantities in linear and rotational motions show striking similarity.

Similar Quantity	Linear	Rotational
Mass	Mass $M$	Moment of Inertia $I = \int r^2 dm$
Length of motion	Distance $L$	Angle $q$ (Radian)
Speed	$v = \frac{dr}{dt}$	$w = \frac{dq}{dt}$
Acceleration	$a = \frac{dv}{dt}$	$\alpha = \frac{dw}{dt}$
Force	Force $F = ma$	Torque $\tau = I\alpha$
Work	Work $W = \int_{x_i}^{x_f} F dx$	Work $W = \int_{q_i}^{q_f} \tau dq$
Power	$P = \vec{F} \cdot \vec{v}$	$P = \tau w$
Momentum	$\vec{p} = m \vec{v}$	$\vec{L} = I \vec{w}$
Kinetic Energy	Kinetic $K = \frac{1}{2} m v^2$	Rotational $K_R = \frac{1}{2} I w^2$

