PHYS 5326 – Lecture #17

Monday, Mar. 24, 2003 Dr. Jae Yu

Mid-term problem reviewMass Terms in LagrangiansSpontaneous Symmetry Breaking

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- Use large number of protons on target to produce many secondary hadrons (π, K, D, etc) and focus as many of them as possible
- Let π and K decay in-flight for v_{μ} beam in the decay pipe $\pi \rightarrow \mu + v_{\mu}$ (99.99%), K $\rightarrow \mu + v_{\mu}$ (63.5%)
- Let the beam go through shield and dirt to filter out μ and remaining hadrons, except for ν
 - Dominated by ν_{μ}

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2. How can we select sign of neutrinos?

- Neutrinos are electrically neutral
- Need to select the charge of the secondary hadrons from the proton interaction on target
- Sets of Dipoles are used to select desired charges of the secondary hadrons



3. How can there be wrong sign of neutrinos in a sign selected beam?

- Interaction of correct sign secondary hadrons with beamline elements, including dump and shields
 - Act as if a fixed target is hit by hadron beam
- Back-scatter of unused protons into the beamline
- CP violating neutrino oscillations

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5. Structure Functions and PDF's

SF is the description of the collection of point-like particles that forms nucleons while PDF's provide momentum distributions of individual partons within the collection.

Assuming parton model, n-N cross section can be rewritten in $\frac{d^2 \mathbf{s}^{n_T}}{dx dy} = \frac{G_F^2 x s}{\mathbf{p} \left(1 + Q^2 / M_W^2\right)^2} \left[q^{n_T}(x) + (1 - y^2)\overline{q}^{n_T} + 2(1 - y)k^{n_T}(x)\right] \frac{\mathsf{Spin 0}}{\mathsf{partons}}$ $\frac{d^2 \mathbf{s}^{\bar{\mathbf{n}}T}}{dxdy} = \frac{G_F^2 x s}{\mathbf{p} \left(1 + Q^2 / M_W^2\right)^2} \left[\frac{-\bar{\mathbf{n}}T}{q}(x) + (1 - y^2)q^{\bar{\mathbf{n}}T} + 2(1 - y)k^{\bar{\mathbf{n}}T}(x)\right]$ $2xF_1^{\boldsymbol{n}(\bar{\boldsymbol{n}})_T} = 2\left[xq^{\boldsymbol{n}(\bar{\boldsymbol{n}})_T}(x) + xq^{-\boldsymbol{n}(\bar{\boldsymbol{n}})_T}(x)\right]$ Comparing the partonneutrino to proton- $F_{2}^{n(\bar{n})T} = 2 \left[xq^{n(\bar{n})T}(x) + xq^{-n(\bar{n})T}(x) + 2xk^{n(\bar{n})T} \right]$ neutrino SF and PDF's are related as $xF_{3}^{n(\bar{n})T} \neq 2 xq^{n(\bar{n})T}(x) - xq^{-n(\bar{n})T}(x)$ Parity violating PHYS 5326, Spring 2003 Monday, Mar. 6 components Jae Yu If no spin 0, 2xF1=F2

6. PDF Evolution: DGLAP Equations

 The evolution equations by Dokshitzer-Gribov-Lipatov-Altarelli-Parisi provide mechanism to evolve PDF's to any kinematic regime or momentum scale, as a function of momentum transfer scale of the interactions

$$\frac{dG(x,M^2)}{d\ln M^2} = \frac{\boldsymbol{a}_s(\boldsymbol{m}^2)}{2\boldsymbol{p}} \int_x^1 \frac{dy}{y} \left[q^s(y,M^2) P_{Gq}^s\left(\frac{x}{y}\right) + G(y,M^2) P_{GG}\left(\frac{x}{y}\right) \right]$$

 $P_{ij}(x/y)$: Splitting function that is the probability of parton i with momentum y get resolved as parton j with momentum x<y





- Define experimental variable to distinguish NC and CC
- Compare the measured ratio with MC prediction

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8-1. $sin^2 \Theta_W$ Theoretical Uncertainty

 Significant correlated error from CC production of charm quark (m_c) modeled by slow rescaling mechanism



• Suggestion by <u>Paschos-Wolfenstein</u> by separating v and \overline{v} beams:

$$R^{-} = \frac{s_{NC}^{n} - s_{NC}^{\overline{n}}}{s_{CC}^{n} - s_{CC}^{\overline{n}}} = ?^{2} \left(\frac{1}{2} - \sin^{2}?_{W}\right) = \frac{R^{n} - R^{\overline{n}}}{1 - r}$$

→ Reduce charm CC production error by subtracting sea quark contributions

 \rightarrow Only valence u, d, and s contributes while sea quark contributions cancel out

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 \rightarrow Massive quark production through Cabbio suppressed d_v quarks only

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8-2. Experimental Uncertainties from $\nu_{\rm e}$

- Electron neutrinos, $\nu_{e'}$ in the beam fakes NC events from CC interactions
 - If the production cross section is well known, the effect will be smaller but since majority come from <u>neutral K (K_L)</u> <u>whose x-sec is known only to 20%</u>, this is a source of large experimental uncertainty
- Using tilted incident proton beam to eliminate neutral hadrons from the secondary beam.

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9. Neutrino Oscillation & Its importance

- Caused by the fact that there are two different eigenstates for mass and weak flavors
- The weak eigenstates are expressed as a linear combination of mass eigenstates with time phase and mixing angle
- Neutrinos are one of the fundamental constituents in nature
 - Three weak eigenstates based on SM
- Left handed particles and right handed anti-particles only
 - Violates parity \rightarrow Why only neutrinos?
 - Is it because of its masslessness?
- SM based on massless neutrinos
- SM inconsistent

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10. Atmospheric Neutrinos & Their Flux

- Neutrinos resulting from the atmospheric interactions of cosmic ray particles
 - He, p, etc + N $\rightarrow \pi$,K, etc
 - $\pi \not \rightarrow \mu + \nu_{\mu}$

 $\mu \not \rightarrow e + \nu_e + \nu_\mu$

- This reaction gives 2 ν_{μ} and 1 ν_{e}
- Expected flux ratio between ν_{μ} and ν_{e} is 2 to 1
- Give a predicted ratio of

$$\left(\frac{N_{n_e}}{N_{n_m}}\right) \approx \frac{1}{2}$$

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11. Importance of Zenith Angle

- The Zenith angle represents the different distance the neutrinos traveled through the earth
- The dependence to the angle is a direct proof of the oscillation probability $P(\mathbf{n}_m \to \mathbf{n}_e) = \sin^2 2q \sin^2 \left(\frac{1.27\Delta m^2 L}{L}\right)$





Super-K Atmospheric Neutrino Results

12. Local Gauge Invariance

Physical meaning of local gauge invariance is the preservation of energy-momentum conservation and the physical law governing the interactions.

Requiring local gauge invariance forces the *L* to accept new vector fields which are massless to preserve the invariance and introduces interactions between the vector field and the field currents.

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Homework

- Presentation of detailed cuts used in data selection
 - Next Monday, Mar. 31

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