PHYS 3446 – Lecture #2

Monday, Jan. 24, 2005 Dr. **Jae** Yu

- 1. Introduction
- 2. History on Atomic Models
- 3. Rutherford Scattering
- 4. Rutherford Scattering with Coulomb force
- 5. Scattering Cross Section
- 6. Measurement of Cross Sections



Announcements

- World scientific acquired 30 copies of the book
- You might want to make sure that the book store knows you want a copy
- Until the book comes, the author generously allowed us to copy the relevant part of the book and use it.
- I have five subscribed the distribution list. You still have time for 5 extra credit points...
- Class assignments:

Most importantly...

Mon

- $W \rightarrow e_{v+X}$: James, Carlos, Justin and Elisha
- W \rightarrow $\mu\nu$ +X: Jeremy, Jay, Jason and Jim
- $Z \rightarrow ee+X$: Casey, David and Mathew
- $Z \rightarrow \mu \mu + X$: John, Jacom and Sabine
- Your accounts will soon be created and a new machine for your login will be available at the end of the week.

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We should have fun in the class!

Why do Physics?

Exp. **•** To understand nature through experimental observations and measurements (**Research**) Theory Establish limited number of fundamental laws, usually with mathematical expressions Predict the nature's course \Rightarrow Theory and Experiment work hand-in-hand \Rightarrow Theory works generally under restricted conditions \Rightarrow Discrepancies between experimental measurements and theory are good for improvements \Rightarrow Improves our everyday lives, though some laws can

take a while till we see amongst us



Structure of Matter



Theory for Microscopic Scale, Quantum Mechanics

- Since we deal with extremely small objects, it is difficult to explain the phenomena with classical mechanics and Electro-magnetism
- The study of atomic structure, thus, led us to quantum mechanics → Extremely successful
 - Long range EM force is responsible for holding atom together
 - EM force is sufficiently weak so that the properties of atoms can be estimated reliably based on perturbative QM calculations
- However, when we step into nucleus regime, the simple Coulomb force does not work since the force in nucleus holds positively charged particles together
- The known forces in nature
 - Strong ~ 1
 - Electro-magnetic ~ 10⁻²
 - Weak ~ 10⁻⁵
 - Gravitational ~ 10⁻³⁸



Evolution of Atomic Models

- 1897: J.J. Thompson Discovered electrons
- 1904: J.J. Thompson Proposed a "plum pudding" model of atoms → Negatively charged electrons embedded in a uniformly distributed positive charge





 1911: Geiger and Marsden with Rutherford perform a scattering experiment with alpha particles shot on a thin gold foil



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- 1912: Rutherford proposes atomic model with a positively charged core surrounded by electrons
 - Rutherford postulated a heavy nucleus with electrons circulating it like planets around the sun
 - Deficiency of instability







- 1913: Neils Bohr proposed a quantified electron orbit
 - Electrons can only transition to pre-defined orbits





 1926: Schrodinger proposed an electron cloud model based on quantum mechanics





Rutherford Scattering

- A fixed target experiment with alpha particle as projectile shot on thin gold foil
 - Alpha particle's energy is low → Speed is well below 0.1c (non-relativistic)
- An elastic scattering of the particles
- What are the conserved quantities in an elastic scattering?
 - Momentum
 - Kinetic Energy





$$v_0^2 = v_\alpha^2 + \frac{m_t}{m_\alpha} v_t^2$$

From these two, we obtain $v_t^2 \left(1 - \frac{m_t}{m_\alpha}\right) = 2\vec{v}_\alpha \cdot \vec{v}_t$

•



The Analysis

$$v_t^2 \left(1 - \frac{m_t}{m_\alpha} \right) = 2 \vec{v}_\alpha \cdot \vec{v}_t$$

- If $m_t << m_{\alpha'}$
 - left-hand side become positive
 - v_{α} and v_{t} must be the same direction
 - Using the actual masses

 $m_e \approx 0.5 MeV/c^2$

 $m_{\alpha} \approx 4 \times 10^3 MeV / c^2$

- If $m_t = m_e$, then $m_t/m_{\alpha} \sim 10^{-4}$. $\rightarrow v_{\alpha} \approx v_0$
- Thus, $p_e/p_{\alpha 0}$ <10⁻⁴.
- Change of momentum of alpha particle is negligible



The Analysis

$$v_t^2 \left(1 - \frac{m_t}{m_\alpha} \right) = 2 \vec{v}_\alpha \cdot \vec{v}_t$$

- If $m_t >> m_{\alpha'}$
 - left-hand side become negative
 - v_{α} and v_{t} is opposite direction
 - Using the actual masses $m_t \approx m_{Au} \approx 2 \times 10^5 MeV / c^2$

 $m_{\alpha} \approx 4 \times 10^3 MeV / c^2$

- If $m_t = m_e$, then $m_t/m_{\alpha} \sim 50$. $\rightarrow v_{\alpha} \approx v_0$
- Thus, $p_e/p_{\alpha 0} \sim 2p_{\alpha 0}$.
- Change of momentum of alpha particle is large.



- We did not take into account electro-magnetic force between the alpha particle and the atom
- Coulomb force is a central force and thus a conservative • force
- Coulomb potential between particles with Ze and Z'e electrical charge separated by distance r is $V(r) = \frac{ZZ'e^2}{r}$
- Since the total energy is conserved,

$$E = \frac{1}{2}mv_0^2 = \text{constant} > 0 \implies v_0 = \sqrt{\frac{2E}{m}}$$

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From energy relation, we obtain

Since there is no net torque, the angular momentum (I=rxp) is conserved. \rightarrow The magnitude of the angular momentum is *I=mvb*.

$$l = m\sqrt{2E/mb} = b\sqrt{2mE} \Rightarrow b^2 = l^2/2mE$$

- From the definition of angular momentum, we obtain an equation of motion $d\chi/dt = l/mr^2$ From energy conservation, we obtain another equation of motion motion

$$\frac{dr}{dt} = \pm \sqrt{\frac{2}{m}} \left(E - V(r) + \frac{l^2}{2mr^2} \right)$$
Centrifugal barrier

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Centrifugal barrier

Effective potential

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• Rearranging the terms, we obtain

$$\frac{dr}{dt} = -\frac{l}{mrb} \sqrt{r^2 \left(1 - \frac{V(r)}{E}\right) - b^2}$$

• and
$$d\chi = -\frac{bdr}{r\left[r^2\left(1-\frac{V(r)}{E}\right)-b^2\right]^{1/2}}$$

Integrating this from r₀ to infinity gives the angular distribution of the outgoing alpha particle



- What happens at the DCA? dr
 - Kinetic energy reduces to 0. -

$$\left. \frac{dr}{dt} \right|_{r=r_0} = 0$$

- The incident alpha could turn around and accelerate

We can obtain
$$r_0^2 \left(1 - \frac{V(r_0)}{E}\right) - b^2 = 0$$
This allows us to determine DCA for each

- This allows us to determine DCÁ for a given potential and χ_0 .
- Define scattering angle q as the changes in the asymptotic angles of the trajectory, we obtain

$$\theta = \pi - 2\chi_0 = \pi - 2b \int_{r_0}^{\infty} \frac{dr}{r \left[r^2 \left(1 - \frac{V(r)}{E} \right) - b^2 \right]^{1/2}}$$

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• For a Coulomb potential

$$V(r) = \frac{ZZ'e^2}{r}$$

• DCA can be obtained for a given impact parameter b,

$$r_{0} = \frac{ZZ'e^{2}/E}{2} \left(1 + \sqrt{1 + 4b^{2}E^{2}/(ZZ'e^{2})^{2}} \right)$$

• And the angular distribution becomes

$$\theta = \pi - 2b \int_{r_0}^{\infty} \frac{dr}{r \left[r^2 \left(1 - \frac{ZZ'e^2}{rE} \right) - b^2 \right]^{1/2}}$$



Replace the variable 1/r=x, and performing the integration, we obtain

 $\theta = \pi + 2b\cos^{-1}$

$$\left[\frac{1}{\sqrt{1+4b^2E^2/\left(ZZ'e^2\right)^2}}\right]$$

• This can be rewritten

$$\frac{1}{\sqrt{1+4b^2E^2/(ZZ'e^2)^2}} = \cos\left(\frac{\theta-\pi}{2}\right)$$

• Solving this for b, we obtain

$$b = \frac{ZZ'e^2}{2E}\cot\frac{\theta}{2}$$

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$$b = \frac{ZZ'e^2}{2E}\cot\frac{\theta}{2}$$

- From the solution for b, we can learn the following
 - 1. For fixed b. E and Z'
 - The scattering is larger for a larger value of Z.
 - Makes perfect sense since Coulomb potential is stronger with larger Z.
 - Results in larger deflection.
 - 2. For a fixed b, Z and Z'
 - The scattering angle is larger when E is smaller.
 - If particle has low energy, its velocity is smaller
 - Spends more time in the potential, suffering greater deflection
 - 3. For fixed Z, Z', and E
 - The scattering angle is larger for smaller impact parameter b
 - Makes perfect sense also, since as the incident particle is closer to the nucleus, it feels stronger Coulomb force.



Assignments

- Compute the masses of electron and alpha particles in MeV/c².
- 2. Compute the gravitational and the Coulomb forces between two protons separated by 10⁻¹⁰m and compare their strengths
- 3. Drive the following equations in your book:
 - Eq. # 1.3, 1.16, 1.19, 1.25, 1.32
- These assignments are due next Monday, Jan. 31.

