#### PHYS 3446 – Lecture #19

Wednesday, Apr. 13, 2005 Dr. **Jae** Yu

- Parity
  - Determination of Parity
  - Parity Violation
- Time Reversal and Charge Conjugation
- The Standard Model
  - Quarks and Leptons
  - Gauge Bosons
  - Symmetry Breaking and the Higgs particle



#### Announcements

- The final quiz next Wednesday, Apr. 20
  - At 1:05pm, in the class (SH200)
  - Covers: Ch 10 what we cover on Monday, Apr. 18
- Macros for your project analysis ready and released yesterday morning
- Due for your project write up is Friday, April 22
  - How are your analyses coming along?
- Keep in mind the final, final homework due is Apr. 20.
- I still need to see a few more of you for individual semester grade discussion.



# Project root and macro file locations

- All hanging from the directory /home/venkat/PHYS3446/ •
- W events •
  - I-- W-E-Nu
    - -- MakeTMBTreeClasses\_so.C -- RunMC.C -- TMBTree\_bu.C -- TMBTree\_bu.h
  - I-- W-Mu-Nu
    - -- MakeTMBTreeClasses\_so.C -- RunMC.C -- TMBTree\_bu.C -- TMBTree\_bu.h
- Z events •
  - I-- Z-E-E
    - |-- MakeTMBTreeClasses\_so.C |-- RunMC.C |-- TMBTree\_bu.C -- TMBTree\_bu.h
  - ·-- Z-Mu-Mu
    - |-- MakeTMBTreeClasses\_so.C
    - -- RunMC.C
    - -- TMBTree\_bu.C -- TMBTree\_bu.h



#### Output of $Z \rightarrow e + e + X$ macro



### Gauge Fields and Mediators

- To keep local gauge invariance, new particles had to be introduced in gauge theories
  - U(1) gauge introduced a new field (particle) that mediates the electromagnetic force: Photon
  - SU(2) gauge introduces three new fields that mediates weak force
    - Charged current mediator:  $W^{\scriptscriptstyle +}$  and  $W^{\scriptscriptstyle -}$
    - Neutral current: Z<sup>0</sup>
  - SU(3) gauge introduces 8 mediators for the strong force
- Unification of electromagnetic and weak force SU(2)xU(1) introduces a total of four mediators
  - Neutral current: Photon, Z<sup>0</sup>
  - Charged current: W<sup>+</sup> and W<sup>-</sup>



# Parity

The space inversion transformation (mirror image) → Switch right- handed coordinate system to left-handed

$$\begin{pmatrix} c \ t \\ x \\ y \\ z \end{pmatrix} \xrightarrow{P \ a \ r \ i \ t \ y} \begin{pmatrix} c \ t \\ -x \\ -y \\ -z \end{pmatrix}$$

- How is this different than spatial rotation?
  - Rotation is continuous in a given coordinate system
    - Quantum numbers related rotational transformation are continuous
  - Space inversion cannot be obtained through any set of rotational transformation
    - Quantum numbers related to space inversion is discrete

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#### Determination of Parity Quantum Numbers

- How do we find out the intrinsic parity of particles?
  - Use observation of decays and production processes
  - Absolute determination of parity is not possible, just like electrical charge or other quantum numbers.
  - Thus the accepted convention is to assign <u>+1 intrinsic</u> parity to proton, neutron and the  $\Lambda$  hyperon.
    - The parities of other particles are determined relative to these assignments through the analysis of parity conserving interactions involving these particles.



### Parity Determination

- When the parity is conserved, it can restrict decay processes that can take place.
- Consider a parity conserving decay:  $A \rightarrow B+C$ 
  - Conservation of angular momentum requires both sides to have the same total angular momentum J.
  - If B and C are spinless, their relative orbital angular momentum (l) must be the same as J(=I+s).
  - Thus conservation of parity implies that

$$\eta_A = \eta_B \eta_C \left(-1\right)^l = \eta_B \eta_C \left(-1\right)^J$$

- If the decay products have spin zero, for the reaction to take place we must have  $\eta_A = \eta_B \eta_C$  between the intrinsic parities



### Parity Determination

• Therefore, the allowed decays must have

$$0^+ \rightarrow 0^+ + 0^+$$

$$0^+ \rightarrow 0^- + 0^-$$

$$0^{-} \rightarrow 0^{+} + 0^{-}$$

- Where the spin intrinsic parity if particles are expressed as  $\mathsf{J}^\mathsf{P}$
- The following decays are prohibited under parity conservation  $0^+ \rightarrow 0^+ + 0^-$

$$0^{-} \rightarrow 0^{-} + 0^{-}$$

$$0^- \rightarrow 0^+ + 0^+$$



# Example 1, $\pi^-$ parity

• Consider the absorption of low energy  $\pi^{-}$  in deuterium nuclei

 $\pi^- + d \rightarrow n + n$ 

- The conservation of parity would require  $\eta_{\pi^{-}}\eta_{d}(-1)^{l_{i}} = \eta_{n}\eta_{n}(-1)^{l_{f}}$
- Since the intrinsic parity of deuteron is +1, and that of the two neutrons is +1,

$$\eta_{\pi^{-}} = (-1)^{l_f - l_i} = (-1)^{l_f + l_i}$$

• This capture process is known to proceed from an li=0 state, thus we obtain  $\eta_{\pi^-} = (-1)^{l_f}$ 



#### Example 1, $\pi^{-}$ parity, cont'd

 Since spin of the deuteron J<sub>d</sub>=1, only a few possible states are allowed for the final state neutrons

1) 
$$|nn\rangle = |J| = 1$$
,  $s = 1$ ,  $l_f = 0$  or  $2\rangle$   
2)  $|nn\rangle = |J| = 1$ ,  $s = 1$ ,  $l_f = 1\rangle$   
3)  $|nn\rangle = |J| = 1$ ,  $s = 0$ ,  $l_f = 1\rangle$ 

- Since the two neutrons are identical fermions, their overall wave functions must be anti-symmetric due to Pauli's exclusion principle → leaves only (3) as the possible solution
- Making pion a pseudo-scalar w/ intrinsic parity  $\eta_{\pi^-} = -1$



### Parity Violation

- Till the observation of "τ-θ" puzzle in cosmic ray decays late 1950's, parity was thought to be conserved in (symmetry of) all fundamental interactions
- The  $\tau$  and  $\theta$  particles seem to have identical mass, lifetimes, and spin (J=0) but decay differently  $\theta^+ \to \pi^+ + \pi^0$   $\left(\eta_{\theta^+} = \eta_{\pi^+} \eta_{\pi^0} = (-1)^2 = 1\right)$  $\tau^+ \to \pi^+ + \pi^+ + \pi^ \left(\eta_{\tau^+} = \eta_{\pi^+} \eta_{\pi^-} = (-1)^3 = -1\right)$
- These seem to be identical particles. Then, how could the same particle decay in two different manner, violating parity?



### Parity Violation

- T.D. Lee and C.N. Yang studied all known weak decays and concluded that there were no evidences of parity conservation in weak decays
  - Postulated that weak interactions violate parity
  - See, <u>http://ccreweb.org/documents/parity/parity.html</u> for more interesting readings
- These turned out to be

$$K^{+} \rightarrow \pi^{+} + \pi^{0} \quad (\mathbf{K}_{\pi 2})$$
$$K^{+} \rightarrow \pi^{+} + \pi^{+} + \pi^{-} (\mathbf{K}_{\pi 3})$$



#### Time Reversal

• Invert time from t  $\rightarrow$  - t .

$$t \underline{T} - t$$
  

$$\vec{r} \underline{T} \vec{r}$$
  

$$\vec{p} = m\dot{\vec{r}} \underline{T} - m\dot{\vec{r}} = -\vec{p}$$
  

$$\vec{L} = \vec{r} \times \vec{p} \underline{T} (\vec{r}) \times (-\vec{p}) = -\vec{r} \times \vec{p} = -\vec{L}$$

• How about Newton's equation of motion?

$$m\frac{d^{2}\vec{r}}{dt^{2}} = \vec{F} = \frac{C}{r^{2}}\hat{r} \quad (-1)^{2}\frac{d^{2}\vec{r}}{dt^{2}} = m\frac{d^{2}\vec{r}}{dt^{2}} = \vec{F} = \frac{C}{r^{2}}\hat{r}$$

– Invariant under time reversal

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# Charge Conjugate

• Conversion of charge from Q  $\rightarrow$  - Q .

$$Q \stackrel{\mathbf{C}}{\longrightarrow} -Q$$

$$\vec{E} = c \frac{q}{r^2} \hat{r} \quad \underline{C} \quad c \frac{-q}{r^2} \hat{r} = -\vec{E}$$
$$\vec{B} = cI \int \frac{d\vec{s} \times \hat{r}}{r^2} \quad \underline{C} \quad c(-I) \int \frac{d\vec{s} \times \hat{r}}{r^2} = -\vec{B}$$

• Under this operation, particles become antiparticles

• What happens to the Newton's equation of motion?

$$m\frac{d^{2}\vec{r}}{dt^{2}} = \vec{F} = \frac{C}{r^{2}}\hat{r} \quad \underline{C} \quad m\frac{d^{2}\vec{r}}{dt^{2}} = \frac{q^{2}}{r^{2}}(-1)^{2}\hat{r} = \vec{F}$$

Invariant under charge conjugate
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- Prior to 70's, low mass hadrons are thought to be the fundamental constituents of matter, despite some new particles that seemed to have new flavors
  - Even lightest hadrons, protons and neutrons, show some indication of substructure
    - Such as magnetic moment of the neutron
  - Questioning whether they really are fundamental particles
- In 1964 Gell-Mann and Zweig suggested independently that hadrons can be understood as composite of quark constituents
  - Recall that the quantum number assignments, such as strangeness, were only calculational tools rather than real particles



- In late 60's, Jerome Friedman, Henry Kendall and Rich Taylor designed an experiment with electron beam scattering off of hadrons and deuterium at SLAC (Stanford Linear Accelerator Center)
  - Data could be easily understood if protons and neutrons are composed of point-like objects with charges -1/3e and +2/3e.
  - A point-like electrons scattering off of point-like quark partons inside the nucleons and hadrons
    - Correspond to modern day Rutherford scattering
    - Higher energies of the incident electrons could break apart the target particles, revealing the internal structure



- Elastic scattering at high energies can be described well with the elastic form factors measured at low energies, why?
  - Since the interaction is elastic, they behave as if they are point-like particles
- Inelastic scattering, on the other hand, cannot be since the target is broken apart
  - Inelastic scatterings of electrons with large momentum transfer (q<sup>2</sup>) provides opportunities to probe shorter distances, breaking apart nucleons
  - The fact that the form factor for inelastic scattering at large q<sup>2</sup> is independent of q<sup>2</sup> shows that there are point-like object in a nucleon
- Nucleons contain both quarks and glue particles (gluons) both described by individual characteristic momentum distributions (Parton Distribution Functions)



- By early 70's, it was clear that hadrons are not fundamental point-like objects
- But leptons did not show any evidence of internal structure
  - Event at very high energies they still do not show any structure
  - Can be regarded as elementary particles
- The phenomenological understanding along with observation from electron scattering (Deep Inelastic Scattering, DIS) and the quark model
- Resulted in the Standard Model that can describe three of the four known forces along with quarks, leptons and gauge bosons as the fundamental particles



#### Quarks and Leptons

In SM, there are three families of leptons  $\begin{pmatrix} V_e \\ e^- \end{pmatrix} \qquad \begin{pmatrix} V_\mu \\ \mu^- \end{pmatrix} \qquad \begin{pmatrix} V_\tau \\ \tau^- \end{pmatrix}$ ➔ Increasing order of lepton masses Convention used in strong isospin symmetry, higher member of multiplet carries higher electrical charge And three families of quark constituents  $\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix} +\frac{2}{3}$ All these fundamental particles are fermions w/ spin  $\frac{1}{2}\hbar$ 



#### Standard Model Elementary Particle Table

• Assumes the following fundamental structure:



• Total of 6 quarks, 6 leptons and 12 force mediators form the entire universe



#### **Quark Content of Mesons**

- Meson spins are measured to be integer.
  - They must consist of an even number of quarks
  - They can be described as bound states of quarks
- Quark compositions of some mesons
  - PionsStrange mesons $\pi^+ = u\overline{d}$  $K^+ = u\overline{s}$  $\pi^- = \overline{u}\overline{d}$  $K^- = \overline{u}\overline{s}$  $\pi^0 = \frac{1}{\sqrt{2}}(u\overline{u} d\overline{d})$  $K^0 = d\overline{s}$  $\overline{K}^0 = \overline{ds}$



### Assignments

1. No homework today!!!

